



Revealing pulsar radio emissions in electron-positron plasmas
by massively parallel Particle-in-cell code simulations

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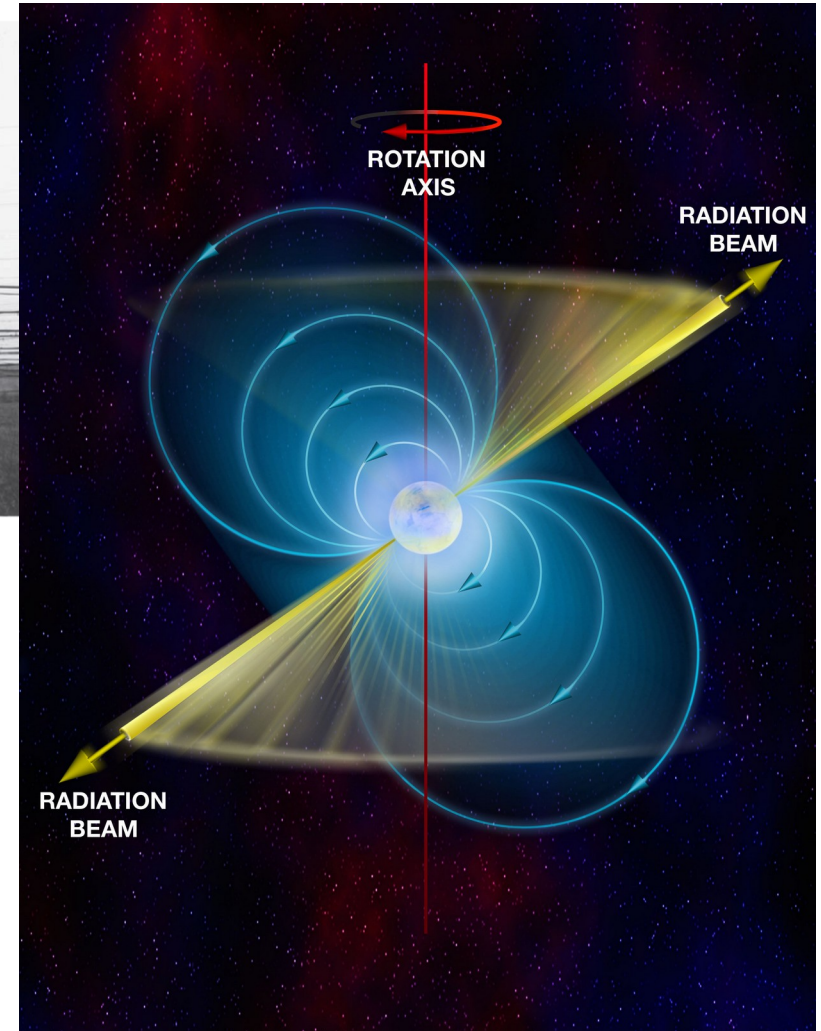
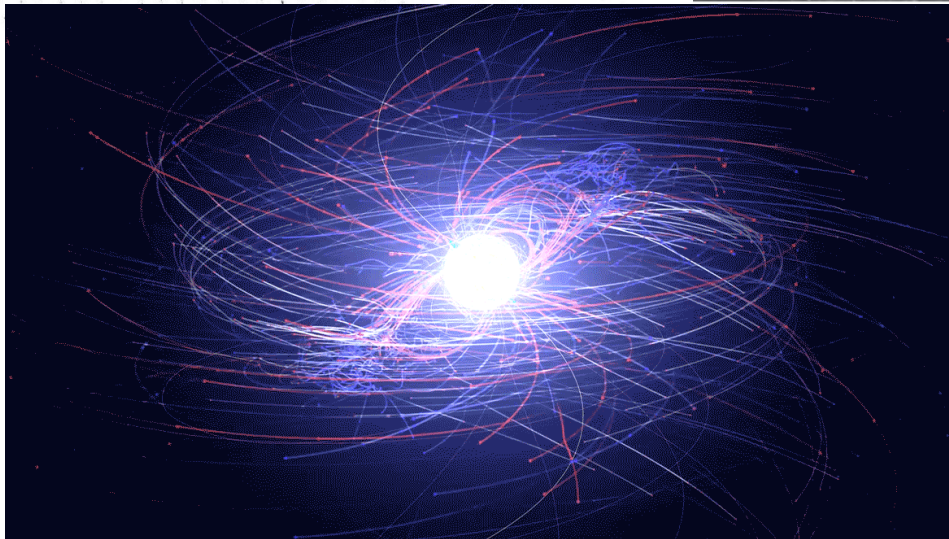
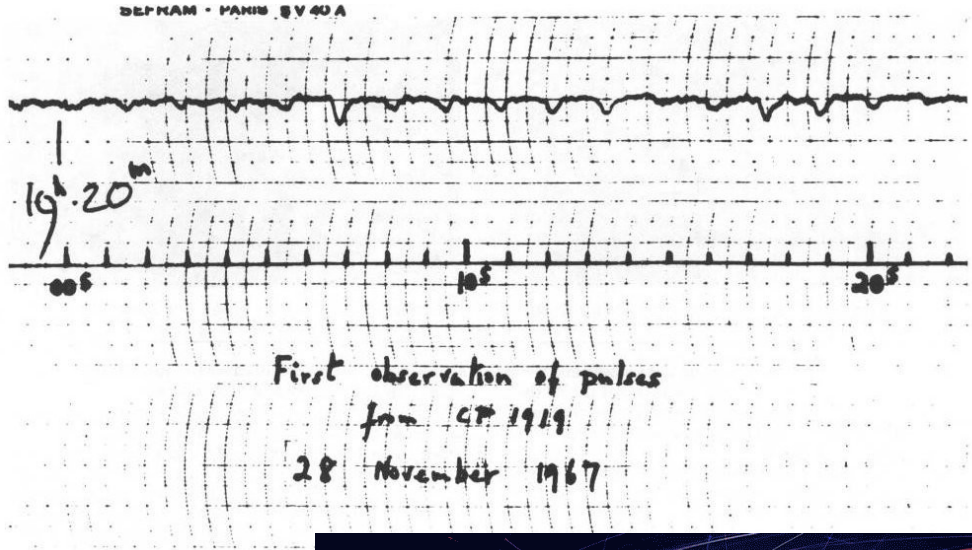
Centre for Astronomy and Astrophysics, Technical University of Berlin

In collaboration with: Patricio Muñoz, Jörg Büchner, Alina Manthei, and Axel Jessner

SuperMUC-NG Status and Results Workshop 2023

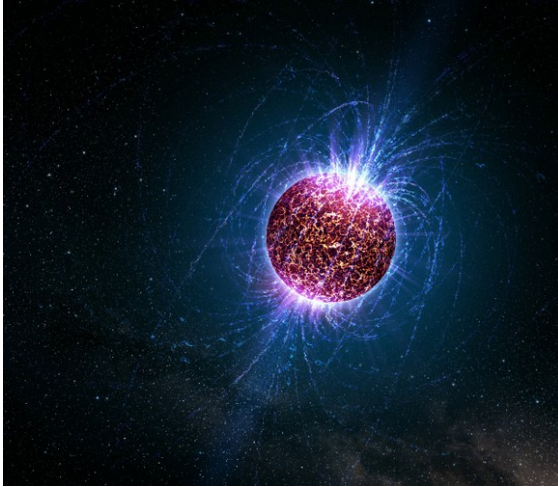
9.5.2023

Pulsar radio emission mechanisms are uncertain



Neutron stars provide insights into a broad range of various astrophysical phenomena

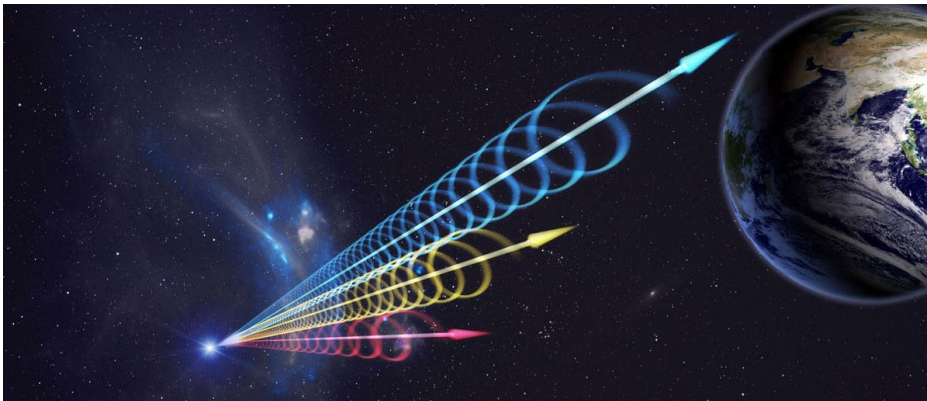
Magnetars



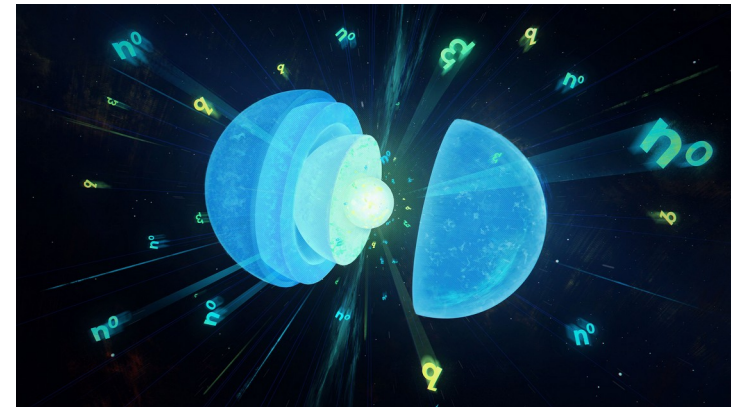
Gravitational waves



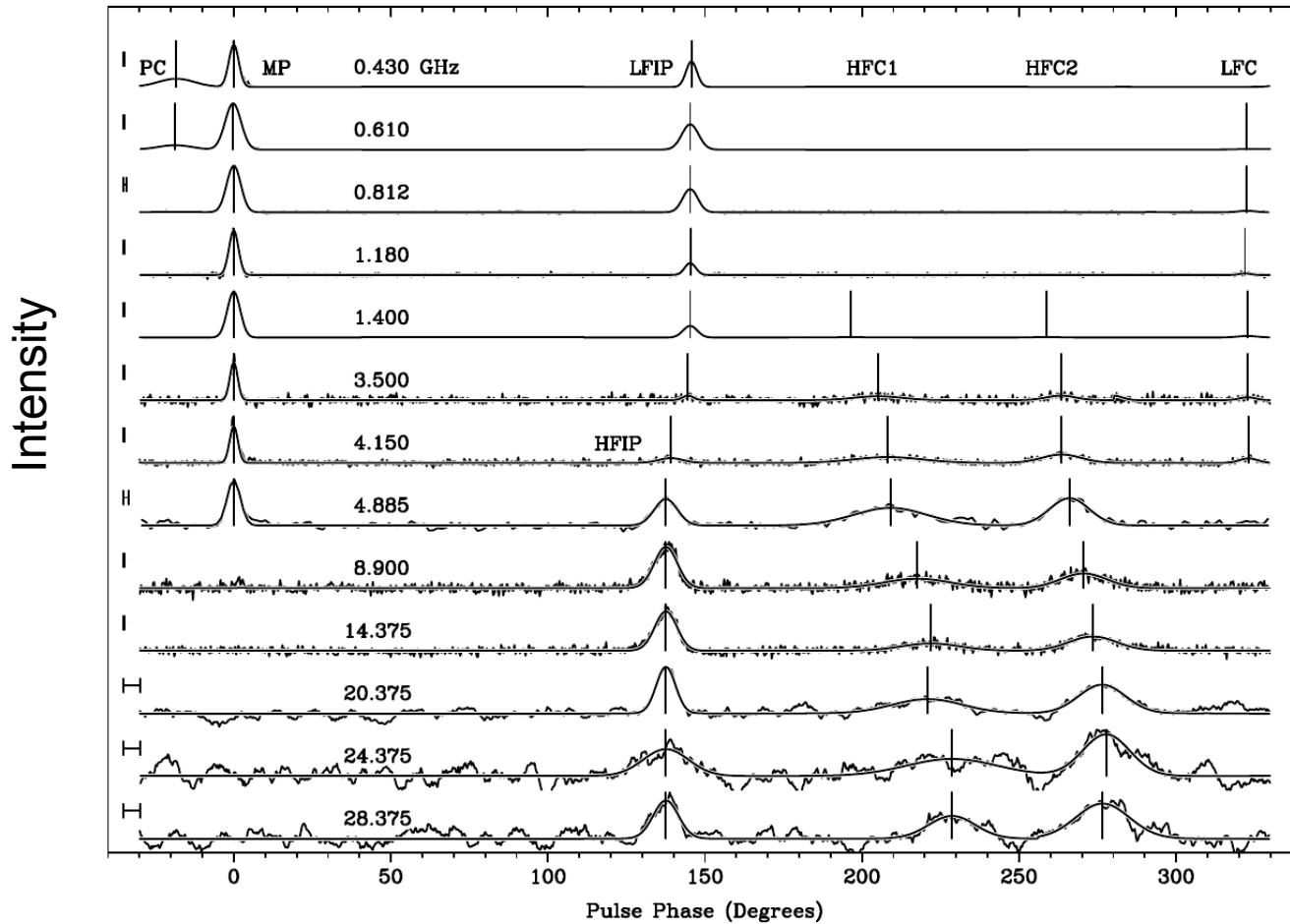
Fast radio bursts



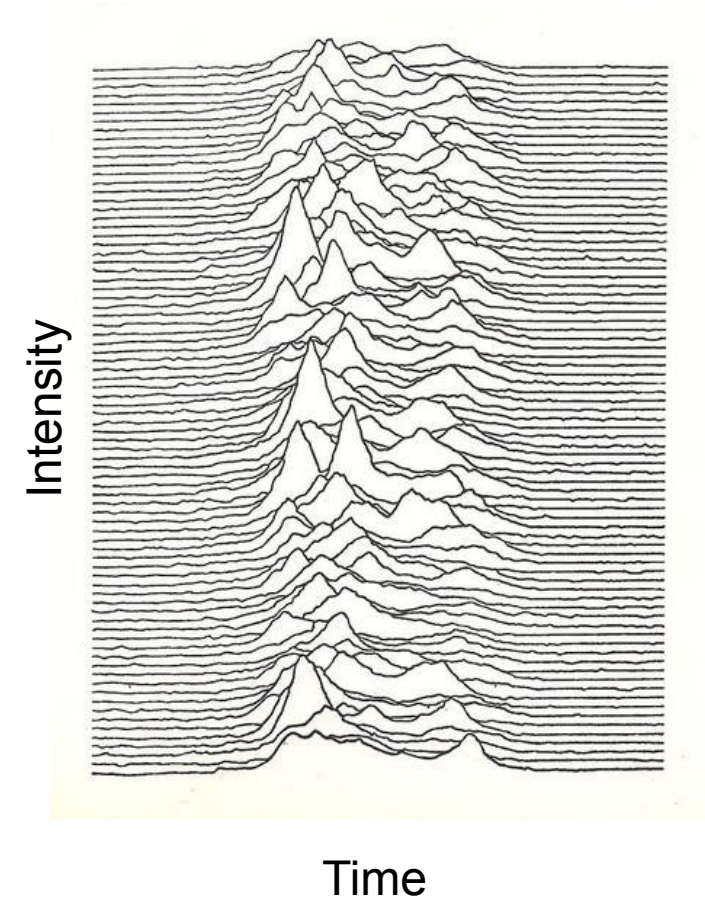
Extreme physical environments



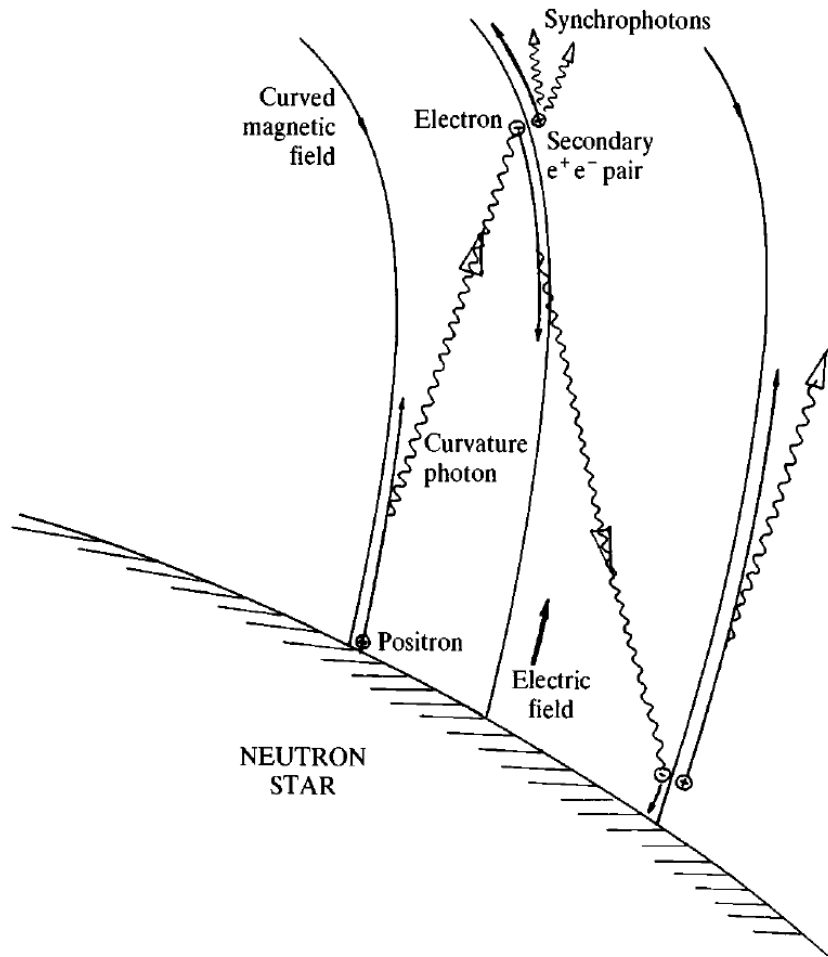
Radio pulses provide wide range of features on various timescales



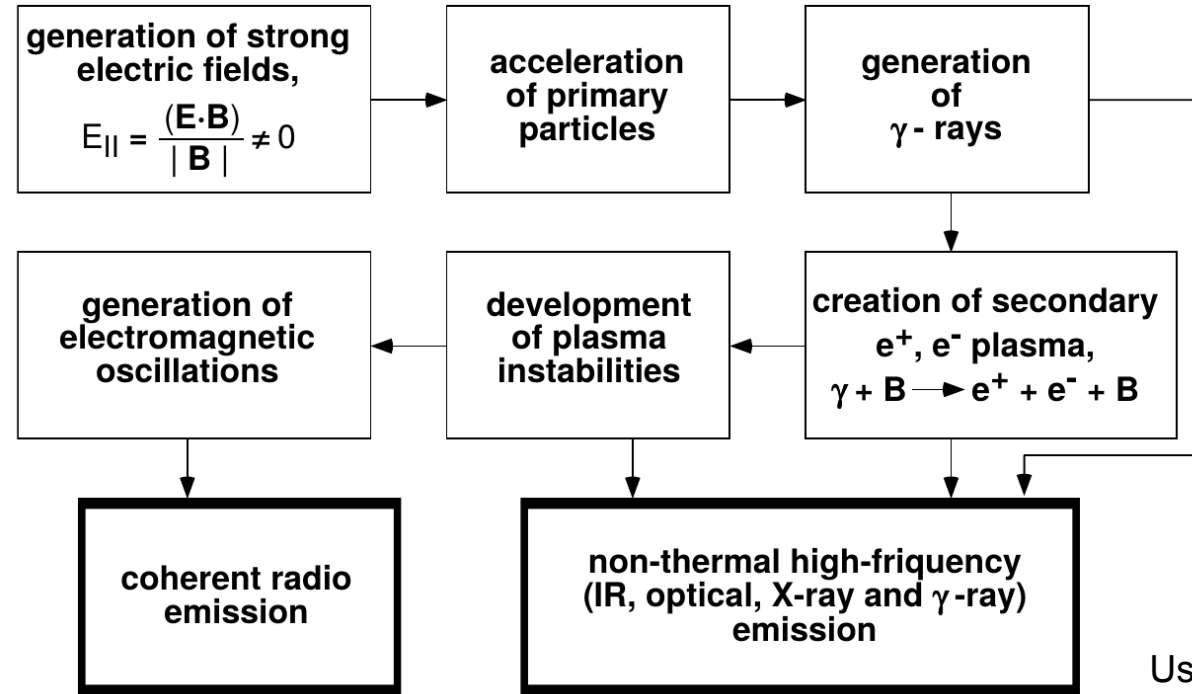
Crab pulsar, Hankins et al. (2015)



Sparking mechanism is a source of electron-positron plasma in gap regions



Beskin et al. (1993)



Usov (2002)

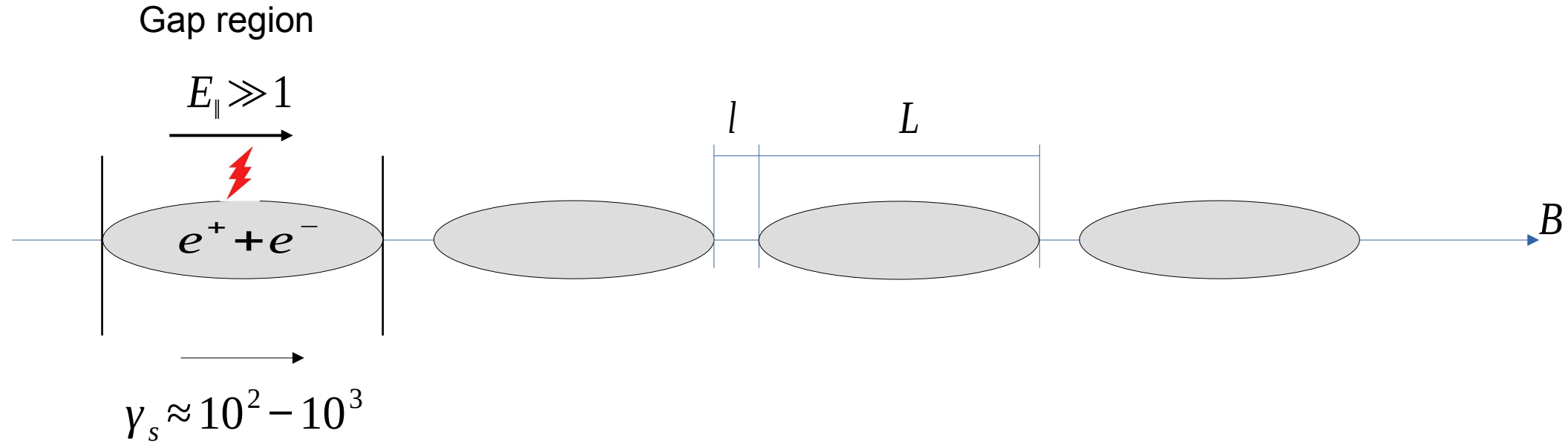
Polar cap region:

$$\gamma + \vec{B} \rightarrow e^+ + e^- + \vec{B}$$

Magnetospheric reconnection:

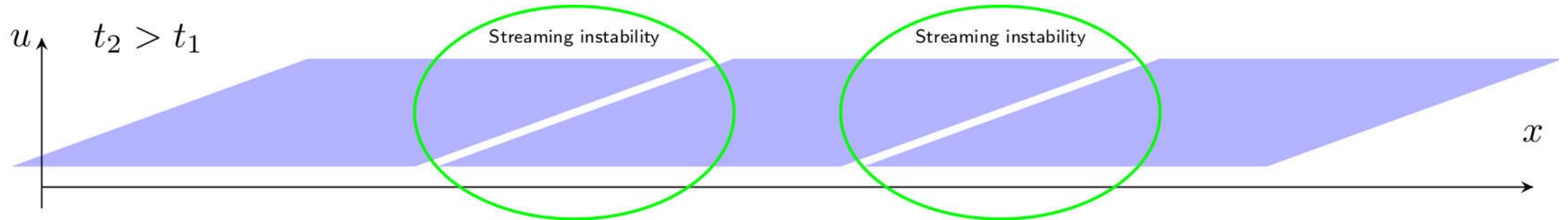
$$\gamma + \gamma \rightarrow e^+ + e^-$$

Train of particle clouds (bunches) outflows from the gap region



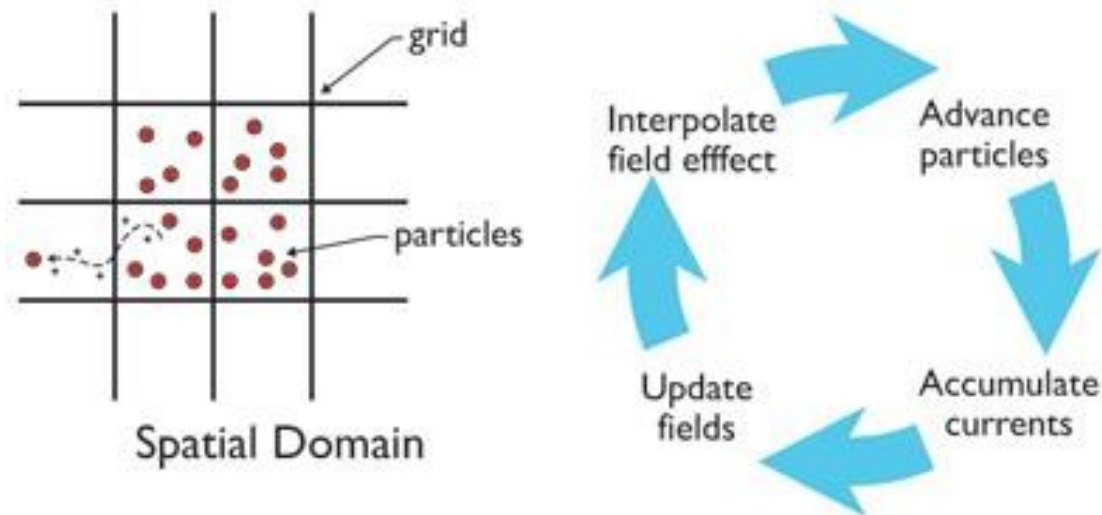
Sturrock (1971)

1.) Streaming instability is produced by bunches overlapping in phase space



Usov & Ursov (1987), Usov (1987), Weatherall (1992), Usov (2002)

Particle-in-cell simulation allows study plasma at kinetic scales



PIC codes:

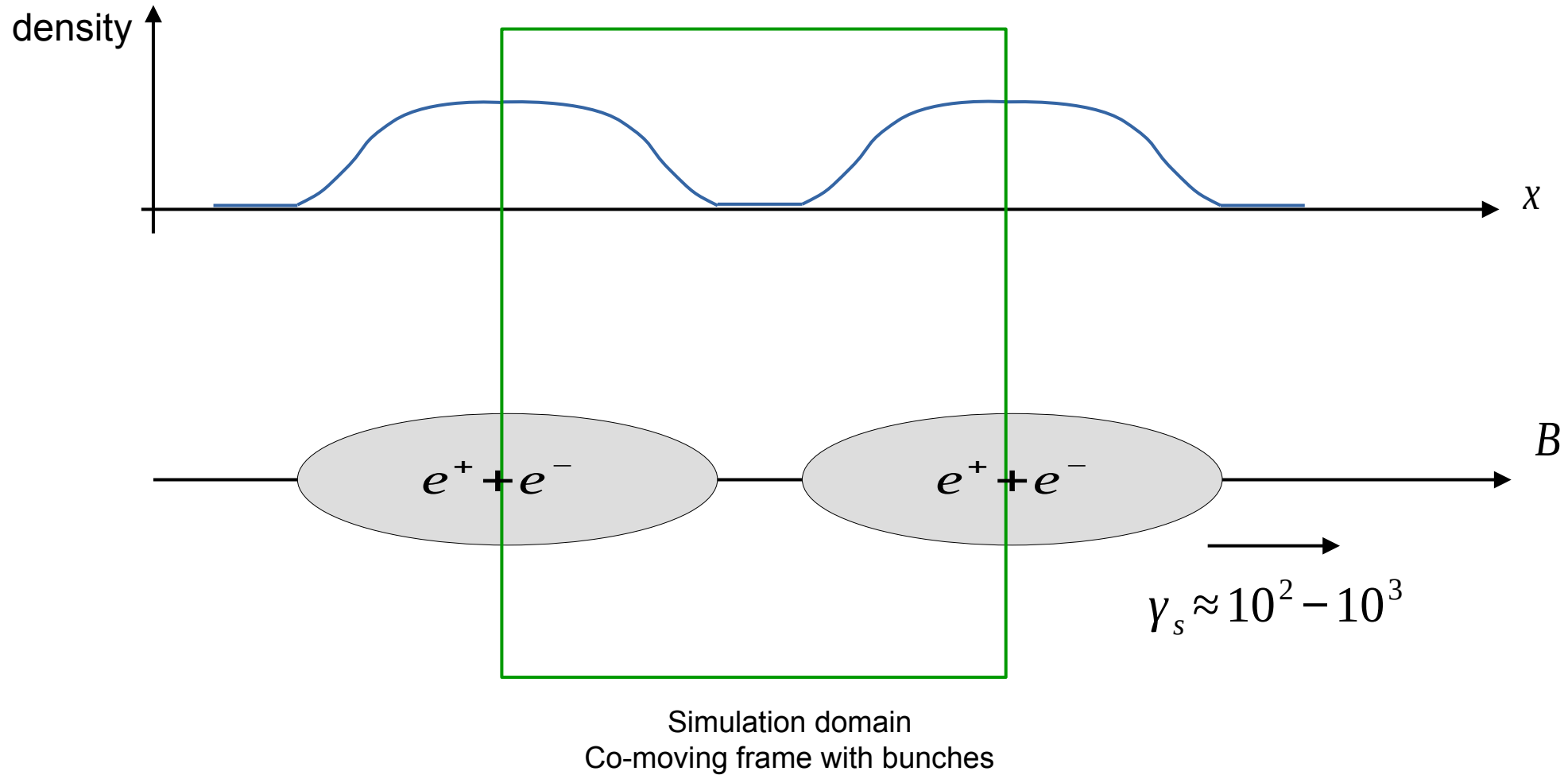
- Electric and magnetic fields on grids
- Velocity distribution sampled into (macro)particles

Numerical solving:

- Vlasov-Boltzmann equation
- Maxwell equation

ACRONYM: Kilian et al. (2012)

Bunch-bunch interaction results

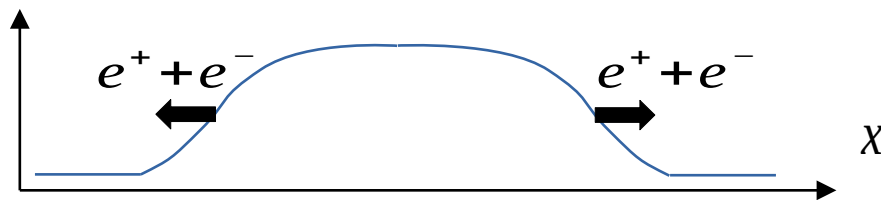


If pairs are created in the gap electric field, they gain initial drift between them

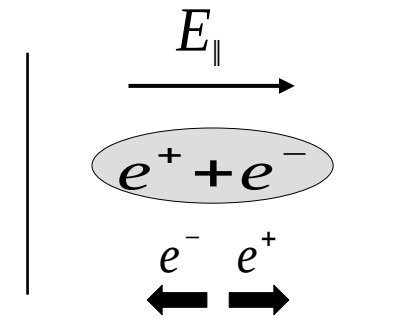
Gap region



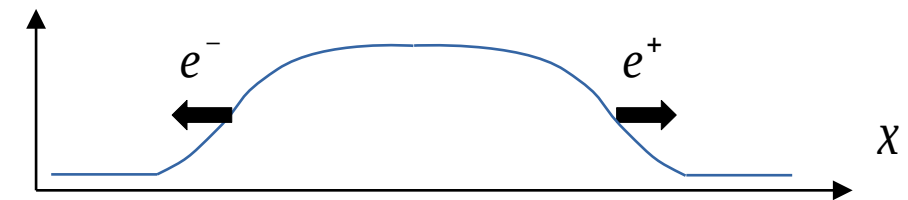
1. Zero drift



Gap region



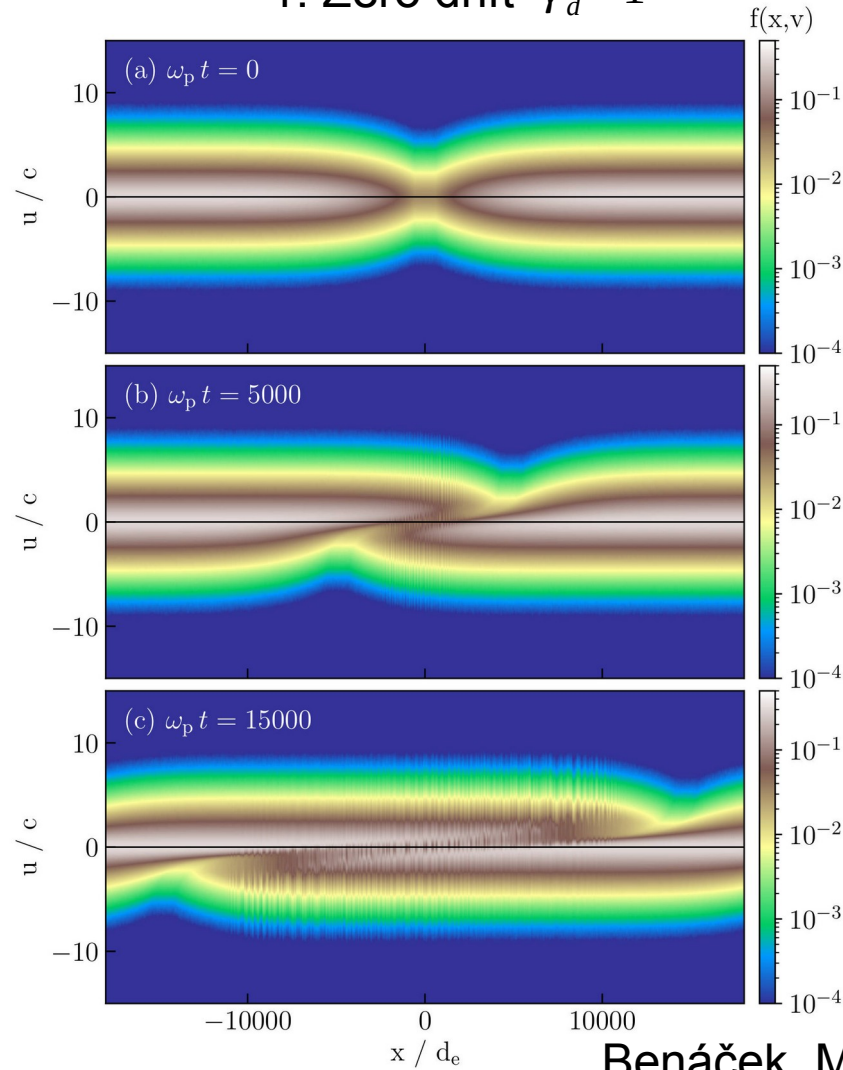
2. Non-zero drift



The bunch expansion completely changes when drift velocity added

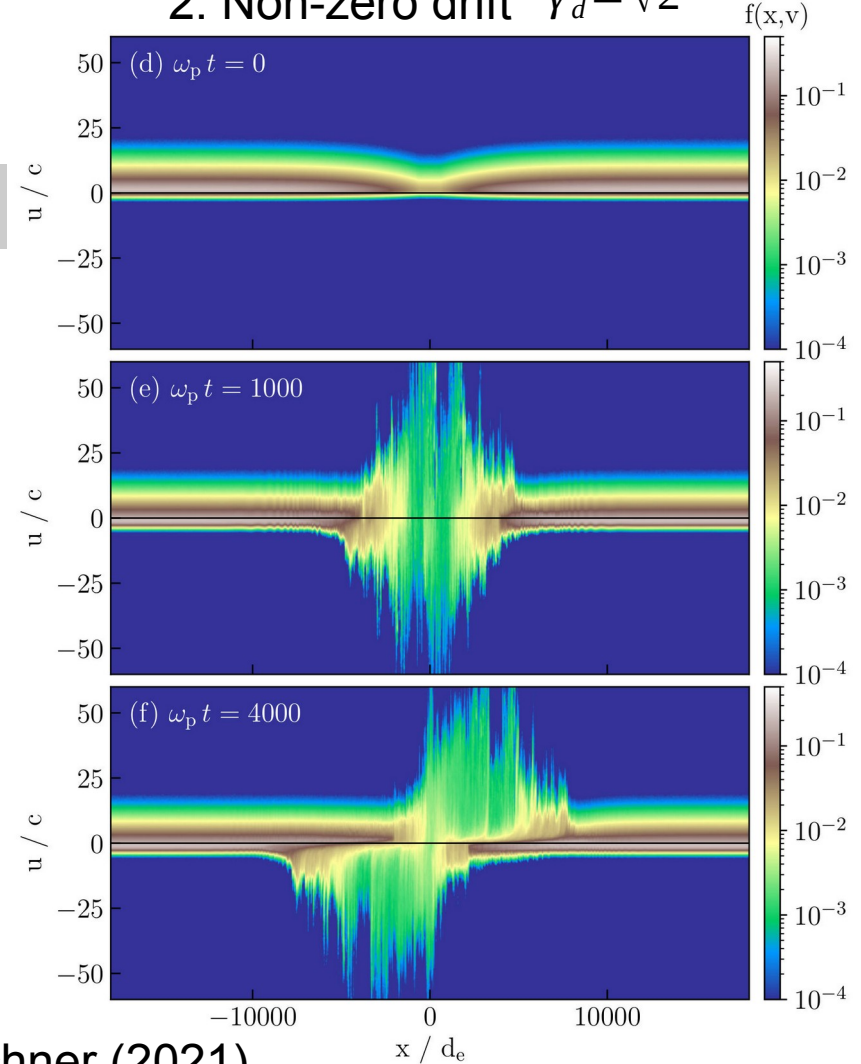
$$u = \frac{p}{m_e}$$

1. Zero drift $\gamma_d = 1$



$$\rho = \frac{mc^2}{k_B T} = 1$$

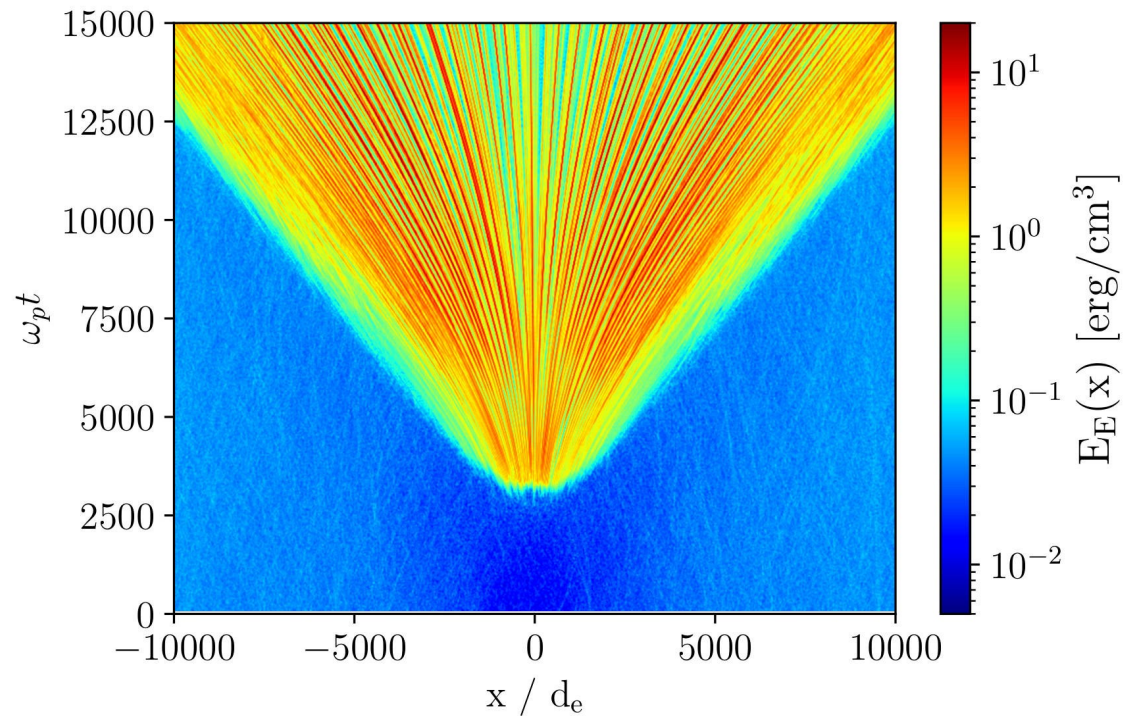
2. Non-zero drift $\gamma_d = \sqrt{2}$



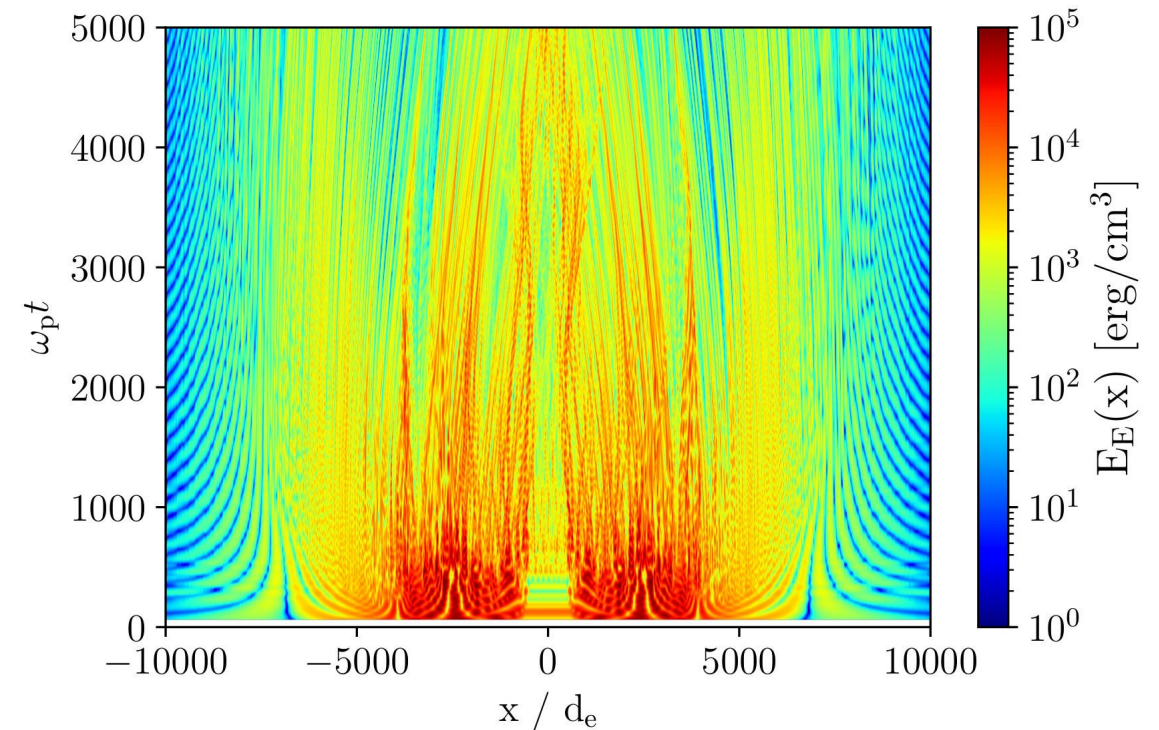
Benáček, Muñoz & Büchner (2021)

The bunch expansion completely changes when drift velocity added

1. Zero drift $\gamma_d=1$

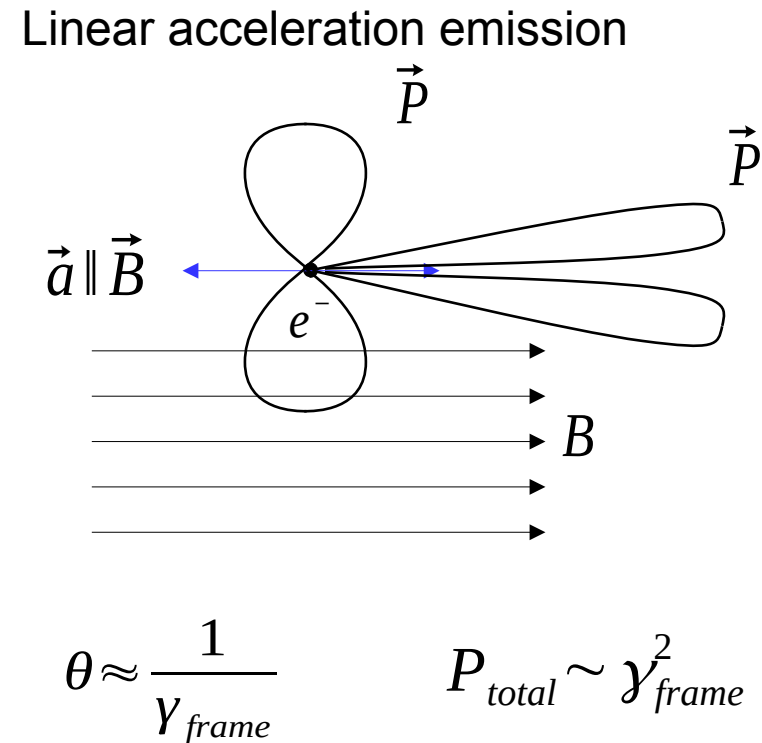
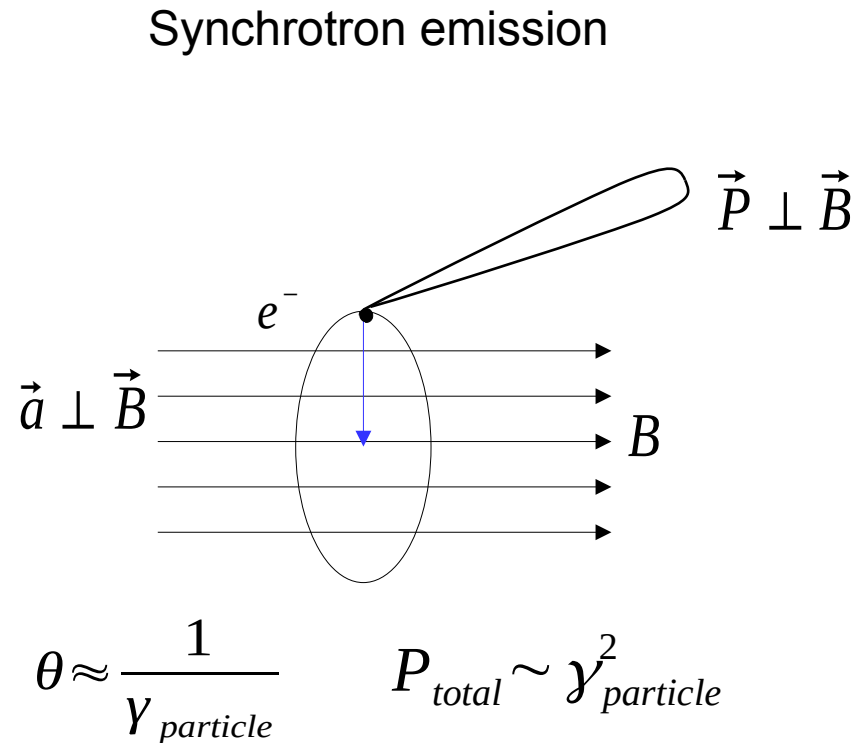


2. Non-zero drift $\gamma_d=\sqrt{2}$



Benáček, Muñoz & Büchner (2021)

Linear acceleration emission vs. Synchrotron emission



Tracking of plasma particles

Using plasma electric current

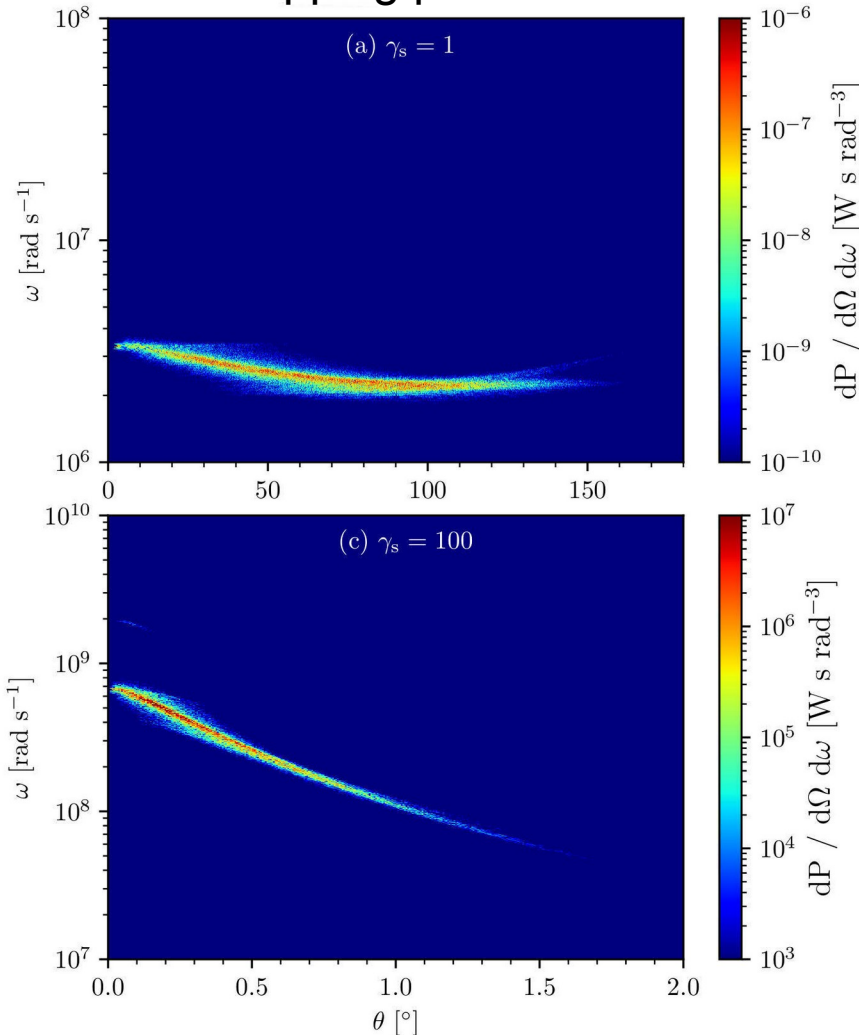
$$\frac{dE}{d\omega d\Omega} = \frac{q^2}{4\pi^2 c} \left| \int_{-\infty}^{+\infty} \frac{\mathbf{n} \times [(\mathbf{n} - \boldsymbol{\beta}) \times \dot{\boldsymbol{\beta}}]}{(1 - \mathbf{n} \cdot \boldsymbol{\beta})^2} e^{i(\omega t - \mathbf{k} \cdot \mathbf{x}(t))} dt \right|^2,$$

$$\frac{dE}{d\omega d\Omega} = \frac{\omega^2}{16\pi^3 \epsilon_0 c^3} | \underline{J(\omega, \mathbf{k})} |^2 \sin^2 \theta,$$

Bunches with nonzero drifts produce significantly stronger radiation

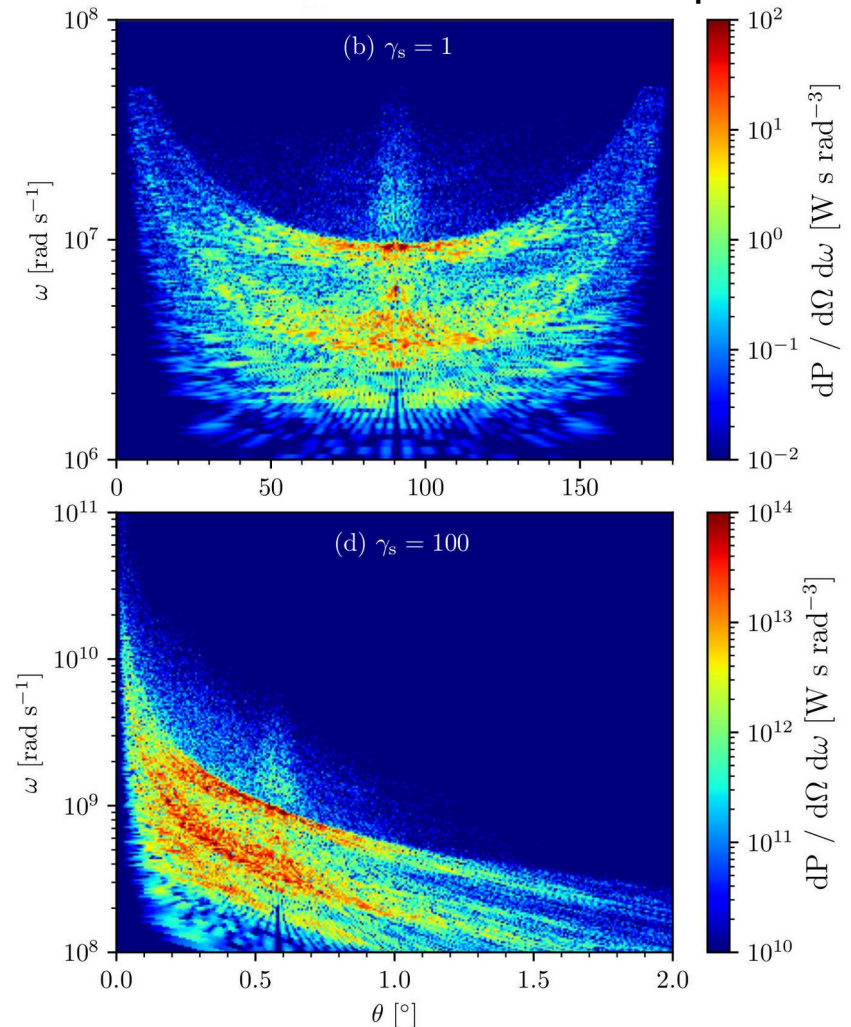
Benáček, Muñoz, Büchner & Jessner (accepted in A&A)

1. Overlapping plasma bunches



$$P_{tot} \approx 3 \times 10^8 \text{ W}$$

2. Bunches constrained from expansion



$$P_{pulsar} \approx 10^{18} - 10^{22} \text{ W}$$

$$P_{tot} \approx 3 \times 10^{16} \text{ W}$$

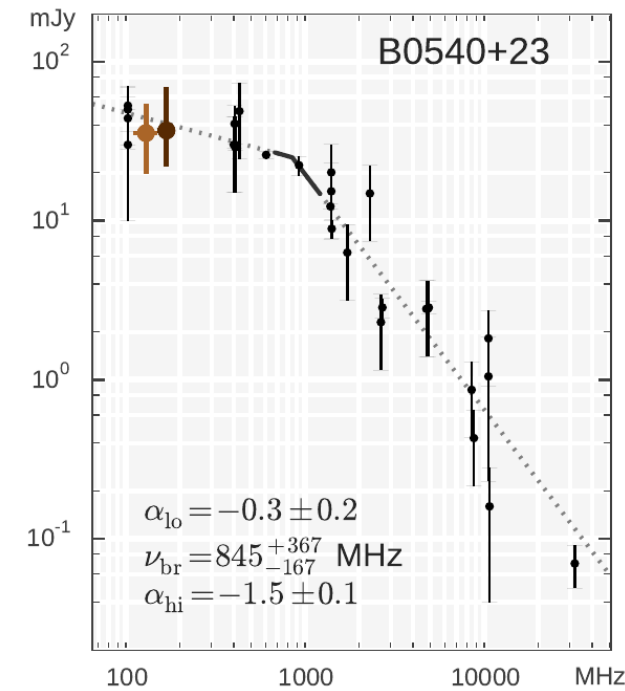
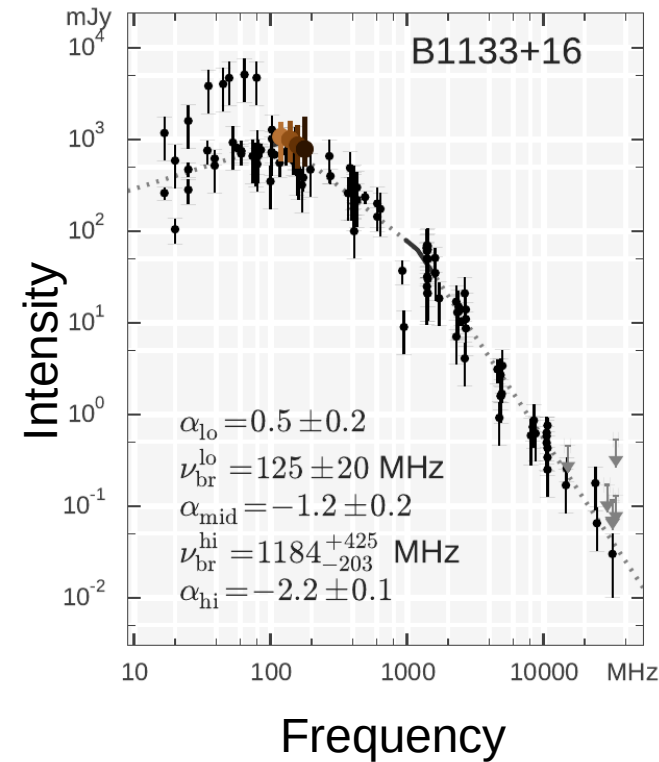
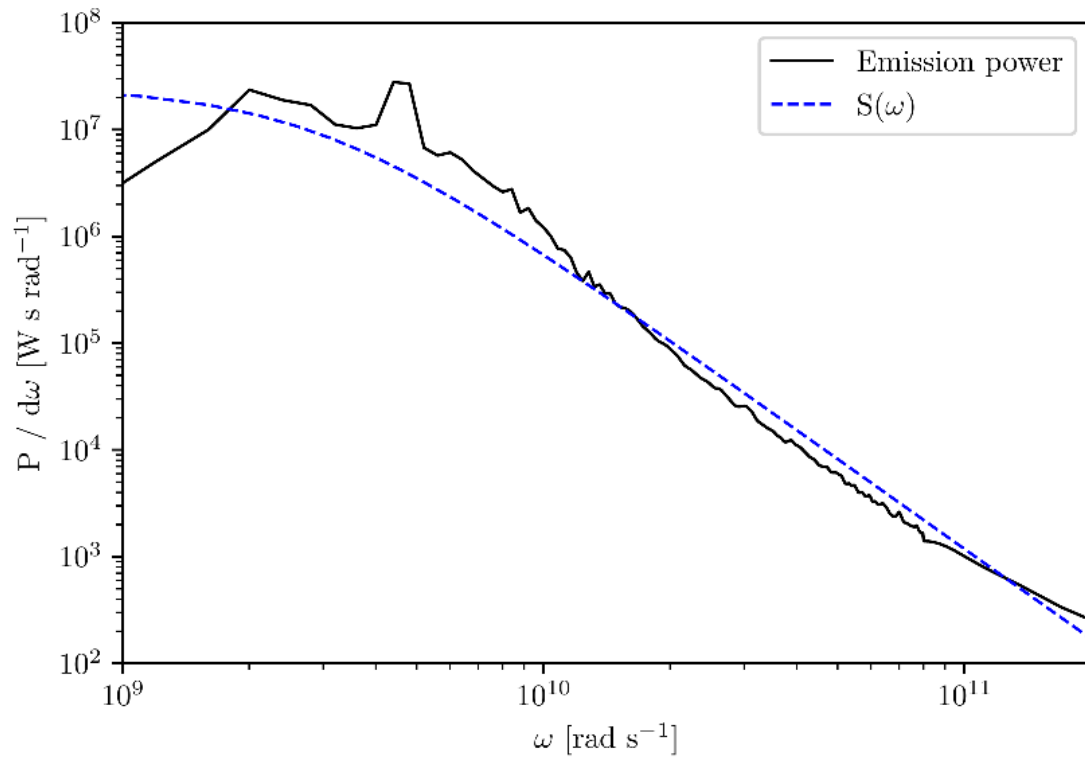
Plasma frame

$$\gamma_{frame} = 1$$

Pulsar frame

$$\gamma_{frame} = 100$$

The frequency profile of emission is similar to those observed

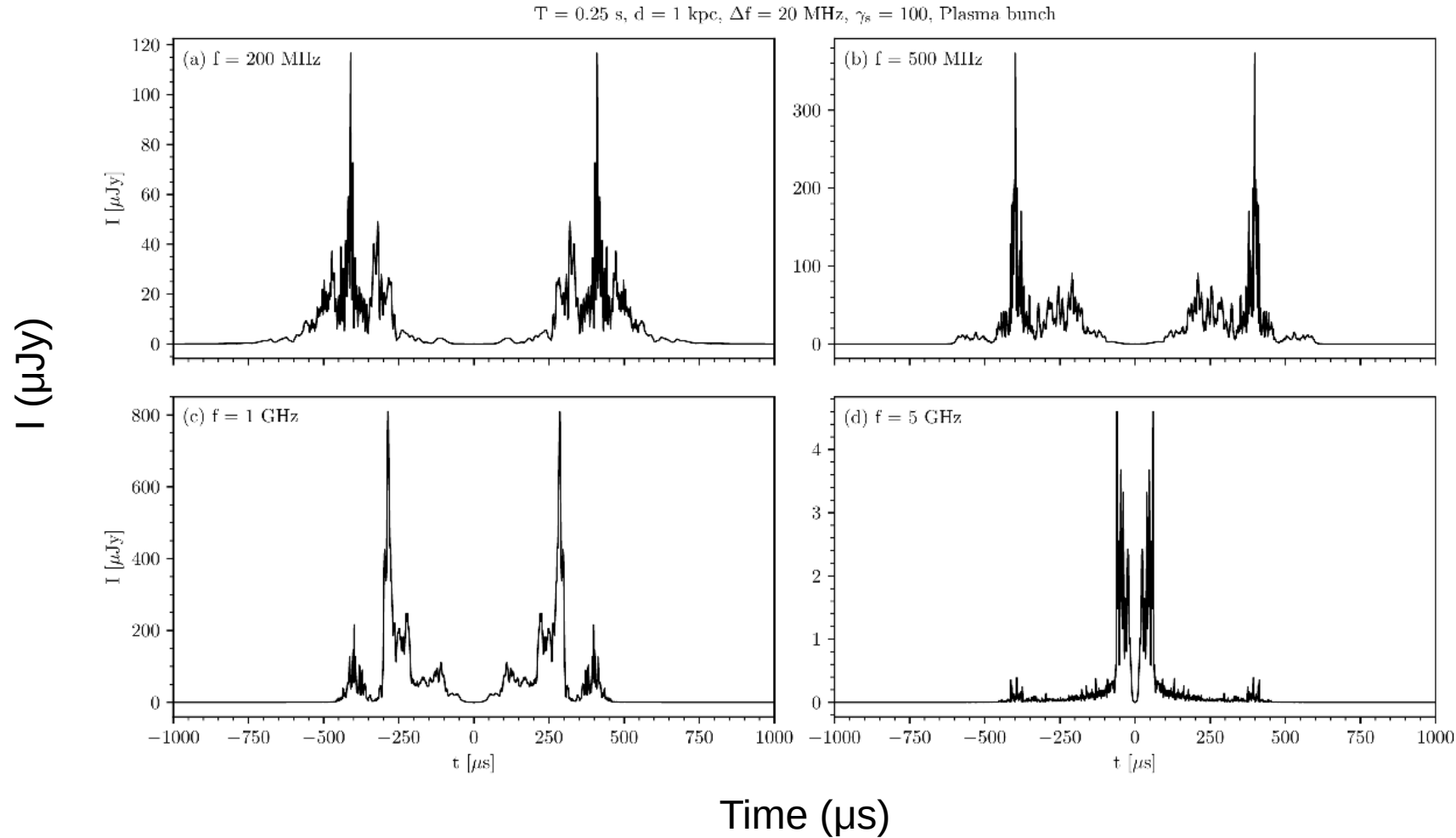


$$S(\omega) = \frac{S_0}{1 - \omega^2 \tau^2}$$

Benáček, Muñoz, Büchner & Jessner (accepted in A&A)

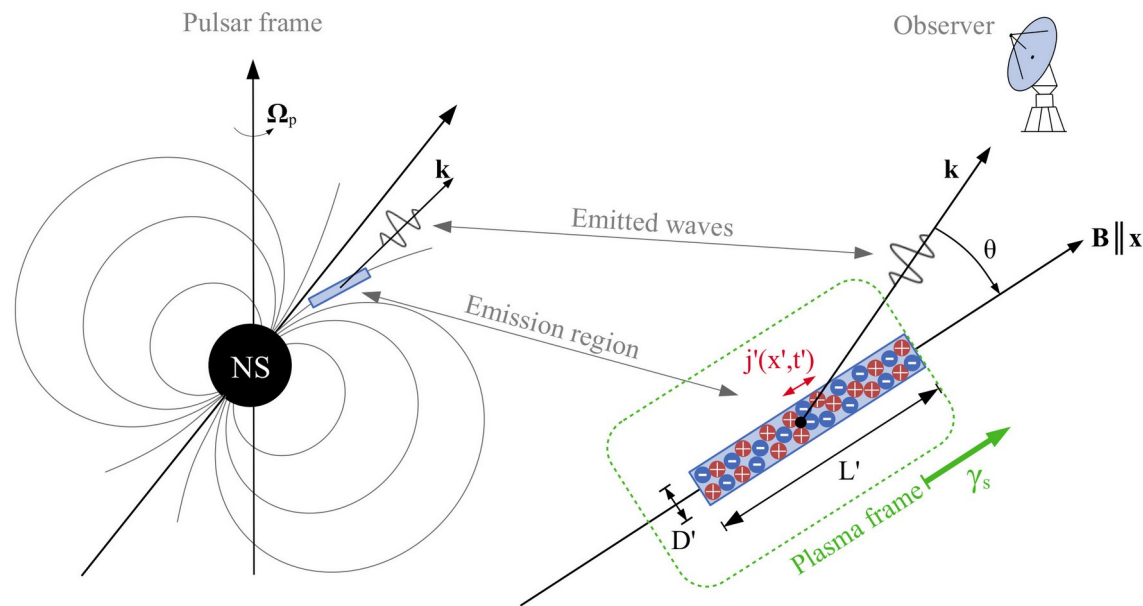
Bilous et al. (2016)

Pulse profiles at specific frequencies



$P_{\text{pulsar}} = 0.25$ s
Distance = 1 kpc
 $\Delta f = 20$ MHz
 $\gamma_{\text{frame}} = 100$

Conclusions



Papers on arXiv:

[2111.05262](https://arxiv.org/abs/2111.05262)

[2106.13525](https://arxiv.org/abs/2106.13525)

[2101.03083](https://arxiv.org/abs/2101.03083)

- Coherent radio emission process of pulsars at kinetic microscales is still uncertain
- Bunch evolution strongly dependent on the initial drift between electrons and positrons
- Linear acceleration emission by bunches might explain the pulsar radiation

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Simulation setup for relativistic streaming instability

Relativistic numerical features

- Relativistic Boris push
- Cole–Kärkkäinen (CK) field solver
- “Weighting with Time dependancy” (WT) shape factor

Simulation setup

- 1D1V code, along mg. field axis
- 720,000 grid cells
- 10^4 PPC ($e^- + e^+$)
- Periodic boundary conditions

Plasma properties

- Maxwell–Jüttner velocity distribution
- Inverse temperature $\rho = mc^2/k_B T = 1$
- Drift between particle species = 1,